



Discrete-Time Rigid Body Pose Estimation Based on Lagrange–d'Alembert Principle

Maulik Bhatt¹ · Srikant Sukumar² · Amit K. Sanyal³

Received: 4 July 2021 / Accepted: 25 August 2022 / Published online: 14 September 2022 © The Author(s), under exclusive licence to Springer Science+Business Media, LLC, part of Springer Nature 2022

Abstract

The problem of rigid body pose estimation is treated in discrete-time via discrete Lagrange-d'Alembert principle and discrete Lyapunov methods. The position and attitude of the rigid body are to be estimated simultaneously with the help of vision and inertial sensors. For the discrete-time estimation of pose, the continuous-time rigid body kinematics equations are discretized appropriately. We approach the pose estimation problem as minimizing the energies stored in the errors of estimated quantities. With the help of measurements obtained through optical sensors, artificial rotational and translation potential energy-like terms have been designed. Similarly, artificial rotational and translation kinetic energy-like terms have been devised using inertial sensor measurements. This allows us to construct a discrete-time Lagrangian as the difference of the kinetic and potential energy-like terms, to which a Lagrange-d'Alembert principle is applied to obtain an optimal pose estimation filter. The dissipation terms in the optimal filter are designed through discrete Lyapunov analysis on a suitably constructed Morse-Lyapunov function, and the overall scheme is proved to be almost globally asymptotically stable. The filtering scheme is simulated using noisy sensor data to verify the theoretical properties.

Keywords Pose estimation · Lagrange–d'Alembert principle · Discrete-time Lyapunov methods

Mathematics Subject Classification 93B27 · 93D20 · 70H30

Communicated by Alain Goriely.

Maulik Bhatt mcbhatt2@illinois.edu; maulikbhatt585@gmail.com

¹ Aerospace Engineering, University of Illinois at Urbana-Champaign, Champaign, IL, USA

² Systems and Control Engineering, Indian Institute of Technology Bombay, Mumbai 400076, India

³ Mechanical and Aerospace Engineering, Syracuse University, Syracuse, NY, USA

1 Introduction

The pose of a rigid body with respect to a frame is a transformation from a body-fixed frame to an inertial frame. The pose encapsulates the position of the center of mass and orientation of the rigid body. Estimation of the pose of a rigid body has various applications in the control of spacecraft, ground vehicles, underwater vehicles for example. Generally, the position and attitude are estimated with the help of onboard inertial sensors coupled with a dynamic model. However, when available, external measurements such as GPS or tracking data of multiple points on the body are also used for pose estimation (Amelin and Miller 2014; Vasconcelos et al. 2008; Vertechy and Castelli 2007). Some techniques combine inertial sensors, vision sensors and external measurements to estimate the rigid body pose. Furthermore, it is common in several applications to operate in GPS denied environments. Therefore, an estimation scheme relying on inertial and vision sensors with proven stability properties and a large domain of attraction is necessary. Additionally, robustness to uncertainties and noise is required.

In recent times, several stable nonlinear estimators evolving on non-Euclidean spaces such as SO(3) or SE(3) have been presented with provably large domain of attraction. A landmark-based nonlinear pose observer is proposed in Vasconcelos et al. (2007) which is almost globally exponentially stable on SE(3). The pose estimation scheme in Rehbinder and Ghosh (2003) uses line-based dynamic vision and inertial sensors to provide a locally convergent attitude observer and subsequently a position estimator. A quaternion-based pose estimator is presented in Filipe et al. (2015) where cost functions based on estimation errors are constructed in discrete-time and minimized to obtain a filtering scheme. The attitude estimation problem based on vector measurements was first proposed as an optimization problem on SO(3) by Wahba (1965). The cost function is known as Wahba's cost function. In Vasconcelos et al. (2010), the authors devise a pose estimator using a Lyapunov function defined as the difference between the estimated and the measured landmark coordinates. For attitude estimation, similar ideas are used in Mahony et al. (2008), Zamani et al. (2013), Izadi and Sanyal (2014), Bhatt et al. (2020b). In Izadi and Sanyal (2016), the authors applied the Lagrange-d'Alembert principle to a Lagrangian constructed through state estimation errors to obtain an optimal filtering scheme for the rigid body pose. However, the work in Izadi and Sanyal (2016) provides a continuous-time pose estimator by applying the (continuous-time) Lagrange-d'Alembert principle to a Lagrangian. The estimator is then discretized for numerical implementation, which voids the theoretical guarantee of asymptotic stability provided by the continuous-time estimator. In this work, we obtain a discrete-time pose estimation scheme by applying the discrete Lagrange-d'Alembert principle on a discrete-time Lagrangian. Furthermore, we also prove guaranteed asymptotic stability by performing the discrete Lyapunov analysis of the system and prove almost global asymptotic stability. Discrete-time observers for *only* attitude with stability properties can be found in Bhatt et al. (2020a, b).

In this paper, we derive an optimal pose estimation scheme by minimizing the "energy" stored in the state estimation errors. A discrete-time Lagrangian has been devised, and the discrete Lagrange–d'Alembert principle from variational mechanics (Marsden and West 2001) is employed to obtain an *optimal* filtering scheme. It is then

proved to be *almost globally asymptotically stable* via discrete Lyapunov analysis. The pose of the rigid body is expressed in SE(3) without employing any local coordinates (such as Euler angles or quaternion) and hence globally non-singular. Furthermore, the estimation scheme presented here relies only on onboard sensor data. We also do not make assumptions on the statistical properties of the measurement noise as is usually the case for Kalman filter-based estimation schemes.

This paper is organized as follows. In Sect. 2, relevant notations are introduced and the procedure to estimate rigid body pose using measurements is explained. The continuous-time rigid body kinematics is discretized in Sect. 3. Section 4 contains the application of variational mechanics to obtain a filter equation for pose estimation. The filter equations obtained in Sect. 4 are proved to be asymptotically stable using the discrete-time Lyapunov method in Sect. 5. Filter equations are numerically verified with realistic measurements (corrupted by bounded noise) in Sect. 6. Finally, Sect. 7 presents the concluding remarks and possible directions future work.

2 Notation and Problem Formulation

2.1 Notation and Preliminaries

We define the trace inner product on $\mathbb{R}^{m \times n}$ as

$$\langle A_1, A_2 \rangle := \operatorname{trace}(A_1^{\mathrm{T}}A_2).$$

The group of orthogonal frame transformations on \mathbb{R}^3 is defined by O(3) := { $Q \in \mathbb{R}^{3\times3} \mid \det(Q) = \pm 1$ }. The special orthogonal group on \mathbb{R}^3 is denoted as SO(3) and defined as SO(3) := { $R \in \mathbb{R}^{3\times3} \mid R^T R = RR^T = I_3$ }. Let there be some $R \in SO(3)$ and $b \in \mathbb{R}^3$. The corresponding Lie algebra is denoted as $\mathfrak{so}(3) := \{M \in \mathbb{R}^{3\times3} \mid M + M^T = 0\}$. The special Euclidean group, SE(3), corresponds to the set of all 4×4 transformation matrices of the form,

$$SE(3) := \left\{ \begin{pmatrix} R & b \\ 0 & 1 \end{pmatrix} \in \mathbb{R}^{4 \times 4} \mid R \in SO(3) \text{ and } b \in \mathbb{R}^3 \right\}.$$

Let $(\cdot)^{\times} : \mathbb{R}^3 \to \mathfrak{so}(3) \subset \mathbb{R}^{3 \times 3}$ be the skew-symmetric matrix cross product operator denoting the vector space isomorphism between \mathbb{R}^3 and $\mathfrak{so}(3)$:

$$v^{\times} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix}^{\times} := \begin{bmatrix} 0 & -v_3 & v_2 \\ v_3 & 0 & -v_1 \\ -v_2 & v_1 & 0 \end{bmatrix}.$$

Further, let $vex(\cdot) : \mathfrak{so}(3) \to \mathbb{R}^3$ be the inverse of $(\cdot)^{\times}$. The map $exp(\cdot) : \mathfrak{so}(3) \to SO(3)$ is defined as

$$\exp\left(M\right) := \sum_{i=0}^{\infty} \frac{1}{k!} M^{k}.$$

We define $Ad : SO(3) \times \mathfrak{so}(3) \to \mathfrak{so}(3)$ as

$$Ad_R \Omega^{\times} := R \Omega^{\times} R^{\mathrm{T}} = (R \Omega)^{\times}.$$

For the remainder of the article, the phrase "consider the time interval $[t_0, T]$ " indicates that the estimation process will be carried out over the time interval $[t_0, T]$ and is divided into N equal subintervals $[t_i, t_{i+1}]$ for i = 0, 1, ..., N with $t_N = T$. The time step size is denoted as, $h := t_{i+1} - t_i$. Further, given a state $z(t), z_i := z(t_i)$. Give a quantity β , β^m denotes its measurement through an onboard sensor.

2.2 Navigation Using Optical and Inertial Sensors

Assume that a rigid body exhibits rotational and translation motion in an environment. The pose estimation of the rigid body implies estimation of the orientation and position of a frame *S*, fixed to the rigid body center of mass with respect to some inertial frame *O*, fixed to the observed environment as shown in Fig. 1. Let $R \in SO(3)$ be the rotation matrix from *S* to *O* and $b \in \mathbb{R}^3$ be the location of the origin of *S* in the frame *O*. We can write the pose $g \in SE(3)$ of the rigid body as,

$$g := \begin{bmatrix} R & b \\ 0 & 1 \end{bmatrix}. \tag{1}$$

If there exists a column vector $\psi = [x \ y \ z]^{T} \in \mathbb{R}^{3}$, then it can be represent as a column vector $\gamma = [x \ y \ z \ 1]^{T}$ in \mathbb{R}^{3} as a subspace of \mathbb{R}^{4} . Furthermore, $g \in SE(3)$ acts on this vector by a combination of rotation and translation as follows: $g\gamma = R\psi + b$.

Assume that there are r^g number of inertial vectors (such as gravity vector, magnetic field vector) whose locations in the frame O are known (denoted as e_j for $j = 1, 2, ..., r^g$) and can be measured in the frame S via inertial sensors (denoted as e_j^s). Furthermore, let there be beacons placed with their position vectors known (r^o in number)in the inertial frame O (denoted as p_j for $j = 1, 2, ..., r^o$). The idea is to measure the locations of these beacons in the vehicle-fixed frame S (denoted a_j) with the help of optical sensors (marked as green in Fig. 1).

It is important to clarify at this stage that at any given discrete-time instant t_i , the number of observed beacons and inertial vectors by the vehicle could be varying. We therefore introduce the notations, r_i^g and r_i^o to denote the corresponding inertial and optical measurements available. It should be evident that $r_i^g \in \{1, 2, ..., r^g\}$ and $r_i^o \in \{1, 2, ..., r^o\}$. We therefore have $\binom{r_i^o}{2}$ unique relative position vectors, which are the vectors connecting any two of these optical beacon measurements. If two or more optical measurements are available, the number of vector measurements that can be used to estimate attitude are $\binom{r_i^o}{2} + r_i^g$. It has to be noted that attitude of the body can be uniquely computed only if $\binom{r_i^o}{2} + r_i^g \ge 2$, . If at least two inertial vector measurements are available then beacon measurements are not required to estimate attitude, however, at least one beacon measurement is necessary for the estimation of relative position. It has been assumed that the velocities of the vehicle can be directly measured.

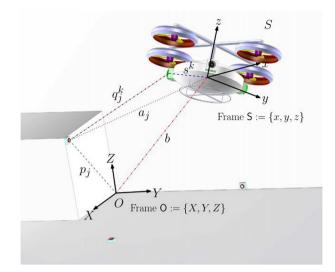


Fig. 1 Inertial landmarks in frame O as observed from vehicle S with optical measurements. *O*—inertial frame, *S*—body-fixed frame, *b*—position of the center of mass of the body, p_j —position of the *j*th beacon in frame *O*, a_j —position of the *j*th beacon in frame *S*, s^k —position of the *k*th optical sensor in frame *S*, q_j^k —range from *k*th optical sensor to *j*th beacon

2.2.1 Pose Measurement Model

Employing the notation from Fig. 1, at the time instant t_i , we obtain

$$p_j = R(q_j^k + s^k) + b = Ra_j + b,$$
 (2)

in the absence of measurement noise. Here $j \in \{1, 2, ..., r_i^o\}$. In the presence of measurement noise, a_i^m can be written as

$$a_j^m = (q_j^k)^m + s^k,$$

Let $\bar{p} = \frac{1}{r_i^o} \sum_{j=1}^{r_i^o} p_j$ be the mean of vectors p_j , and $\bar{a}^m = \frac{1}{r_i^o} \sum_{j=1}^{r_i^o} a_j^m$ be the mean of vectors a_j^m . We obtain the following relation from (2):

$$\bar{a}^m = R^{\mathrm{T}}(\bar{p} - b) + \zeta, \tag{3}$$

where ζ is the additive measurement noise obtained by averaging the measurement noise. As stated in the previous subsection, we obtain $n := \binom{r_0^o}{2}$ relative vectors from optical measurements. They are denoted as $d_j = p_{\lambda} - p_l$ in O and the corresponding vectors in the frame S are denoted as $l_j = a_{\lambda} - a_l$ with $\lambda, l \in \{1, 2, ..., r_i^o\}; \lambda \neq l$. We have

$$d_j = Rl_j \Rightarrow D = RL. \tag{4}$$

🖄 Springer

Putting them in the matrix form as $D = [d_1 \dots d_n]$ and $L = [l_1 \dots l_n] \in \mathbb{R}^{3 \times n}$, we obtain

$$L^m = R^{\mathrm{T}} D + \mathscr{L},\tag{5}$$

where $\mathscr{L} \in \mathbb{R}^{3 \times n}$ consists of the additive noise in the vector measurements made in the body frame *S*.

3 Discretization of Rigid Body Kinematics

Consider the time interval $[t_0, T]$. Let $\Omega \in \mathbb{R}^3$ and $\nu \in \mathbb{R}^3$ be the rotational and translational velocity of the rigid body, respectively, in frame *S*. $R \in SO(3)$ is the rotation matrix from body frame to inertial frame and $b \in \mathbb{R}^3$ is the position of rigid body with respect to frame *O* expressed in frame *S*. The generalized velocity of the rigid body is constructed as $\xi = [\Omega \ \nu]^T$ and the pose of the rigid body is,

$$\mathrm{SE}(3) \ni g = \begin{bmatrix} R & b \\ 0 & 1 \end{bmatrix}.$$

The continuous-time rigid body kinematics are:

$$\dot{R} = R\Omega^{\times}, \quad \dot{b} = R\nu \Rightarrow \dot{g} = g\xi^{\vee},$$

where $\xi^{\vee} := \begin{bmatrix} \Omega^{\times} & \nu \\ 0 & 0 \end{bmatrix}$.

For the discrete-time pose estimation of the rigid body, the continuous-time kinematics are discretized as

$$R_{i+1} = R_i \exp\left(\frac{h}{2}(\Omega_{i+1} + \Omega_i)^{\times}\right), \quad b_{i+1} = b_i + \frac{h}{2}R_{i+1}(\nu_i + \nu_{i+1}).$$
(6)

Therefore, the discrete-time kinematics of the rigid body pose can be expressed as

$$g_{i+1} = g_i \begin{bmatrix} \exp\left(\frac{h}{2}(\Omega_{i+1} + \Omega_i)^{\times}\right) \exp\left(\frac{h}{2}(\Omega_{i+1} + \Omega_i)^{\times}\right)\frac{h}{2}(\nu_i + \nu_{i+1})\\ 0 \end{bmatrix}.$$
 (7)

4 Discrete-Time Estimation of Motion from Measurements

Consider the time interval $[t_0, T]$. Let $\hat{\Omega}_i$ and $\hat{\nu}_i$ be the estimated rotational and translational velocity of the rigid body, respectively, in the frame *S* at time instant t_i . \hat{R}_i is the estimated rotation matrix from *S* to *O* and \hat{b}_i is the estimated position of rigid body with respect to frame *O* expressed in frame *S* at time instant t_i . The generalized estimated velocity of the rigid body is constructed as $\hat{\xi}_i := [\hat{\Omega}_i \ \hat{\nu}_i]^T$. From (7), the

estimated pose and its kinematics can be computed as

$$SE(3) \ni \hat{g}_{i} := \begin{bmatrix} \hat{R}_{i} & \hat{b}_{i} \\ 0 & 1 \end{bmatrix},$$
$$\hat{g}_{i+1} = \hat{g}_{i} \begin{bmatrix} \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \frac{h}{2} \left(\hat{\nu}_{i} + \hat{\nu}_{i+1}\right) \\ 0 & 1 \end{bmatrix}.$$
(8)

The pose estimation error h_i at time instant t_i can be computed as

$$SE(3) \ni \bar{g}_i := g_i \hat{g}_i^{-1} = \begin{bmatrix} Q_i \ b_i - Q_i \hat{b}_i \\ 0 \ 1 \end{bmatrix} = \begin{bmatrix} Q_i \ x_i \\ 0 \ 1 \end{bmatrix}, \tag{9}$$

where $Q_i = R_i \hat{R}_i^{T}$ is the attitude estimation error and $x_i = b_i - Q_i \hat{b}_i$ is the position estimation error. The estimation error in the generalized velocity is denoted as

$$\varphi_i := \varphi(\xi_i^m, \hat{\xi}_i) = \begin{bmatrix} \omega_i \\ \upsilon_i \end{bmatrix} = \xi_i^m - \hat{\xi}_i, \tag{10}$$

where $\omega_i = \Omega_i^m - \hat{\Omega}_i$ is the angular velocity estimation error and $v_i = v_i^m - \hat{v}_i$ is the translational velocity estimation error. Here, Ω_i^m and v_i^m denote the measurements of angular and translational velocities, respectively, at time instant t_i . The discrete-time kinematics of the attitude estimation error and the position estimation error are evaluated as

$$Q_{i+1} = R_{i+1} \hat{R}_{i+1}^{\mathrm{T}} = Q_i \hat{R}_i \exp\left(\frac{h}{2}(\omega_{i+1} + \omega_i)^{\times}\right) \hat{R}_i^{\mathrm{T}},$$
(11)

and

$$\begin{aligned} x_{i+1} &= b_{i+1} - Q_{i+1}\hat{b}_{i+1} \\ &= b_i + R_{i+1}\frac{h}{2}\left(v_i + v_{i+1}\right) - Q_{i+1}\left(\hat{b}_i + \hat{R}_{i+1}\frac{h}{2}\left(\hat{v}_i + \hat{v}_{i+1}\right)\right) \\ &= b_i - Q_{i+1}\hat{b}_i + R_{i+1}\frac{h}{2}\left(v_i + v_{i+1}\right) \\ &= b_i - Q_i\hat{R}_i\exp\left(\frac{h}{2}(\omega_{i+1} + \omega_i)^{\times}\right)\hat{R}_i^{\mathrm{T}}\hat{b}_i + R_{i+1}\frac{h}{2}\left(v_i + v_{i+1}\right).\end{aligned}$$

Approximating $\exp\left(\frac{h}{2}(\omega_{i+1}+\omega_i)^{\times}\right)$ by the first two terms in the expansion as

$$\exp\left(\frac{h}{2}(\omega_{i+1}+\omega_i)^{\times}\right) \approx I + \frac{h}{2}(\omega_{i+1}+\omega_i)^{\times},\tag{12}$$

🖄 Springer

we have.

$$x_{i+1} = b_i - Q_i \hat{b}_i - \frac{h}{2} Q_i \left(\hat{R}_i (\omega_{i+1} + \omega_i) \right)^{\times} + R_{i+1} \frac{h}{2} (v_i + v_{i+1})$$

= $x_i - \frac{h}{2} Q_i \left(\hat{R}_i (\omega_{i+1} + \omega_i) \right)^{\times} + R_{i+1} \frac{h}{2} (v_i + v_{i+1}).$ (13)

It has to be noted that approximation in (12) is accurate for small values of h and may affect the stability results for very high values of h.

4.1 Discrete-time optimal pose estimator based on Lagrange-d'Alembert principle

The error in the attitude estimation is encapsulated by Wahba's cost function (Wahba 1965). Thus, the artificial potential function for rotation estimation error is defined as

$$\mathcal{U}_i^r := \mathcal{U}^r(\hat{g}_i, L_i^m, D_i) = \frac{1}{2} k_p \langle D_i - \hat{R}_i L_i^m, (D_i - \hat{R}_i L_i^m) W_i \rangle, \qquad (14)$$

where D and L^m are as defined in (4) and (5), $W = \text{diag}(w_i) \in \mathbb{R}^{n \times n}$ is a positive definite diagonal matrix of the weight factors for the measured directions, and $k_p > 0$ is a scalar gain. The artificial potential function for translation estimation error is defined as:

$$\mathcal{U}_{i}^{t} := \mathcal{U}^{t}(\hat{g}_{i}, \bar{a}_{i}^{m}, \bar{p}_{i}) = \kappa \|y_{i}\|^{2} := \kappa \|\bar{p}_{i} - \hat{R}_{i}\bar{a}_{i}^{m} - \hat{b}_{i}\|^{2},$$
(15)

where \bar{p} and \bar{a}^m are as per (3) and $\kappa > 0$ is a scalar gain. The total artificial potential energy will be the sum of the artificial rotational and translational potential functions:

$$\mathcal{U}_{i} := \mathcal{U}(\hat{g}_{i}, L_{i}^{m}, D_{i}, \bar{a}_{i}^{m}, \bar{p}_{i}) = \mathcal{U}^{r}(\hat{g}_{i}, L_{i}^{m}, D_{i}) + \mathcal{U}^{t}(\hat{g}_{i}, \bar{a}_{i}^{m}, \bar{p}_{i})$$

$$= \frac{1}{2}k_{p}\langle D_{i} - \hat{R}_{i}L_{i}^{m}, (D_{i} - \hat{R}_{i}L_{i}^{m})W_{i}\rangle$$

$$+ \kappa ||\bar{p}_{i} - \hat{R}_{i}\bar{a}_{i}^{m} - \hat{b}_{i}||^{2}.$$
(16)

We define the artificial kinetic energy of the system as a quadratic in the generalized velocity estimation error:

$$\mathcal{T}_{i} := \mathcal{T}\left(\varphi(\xi_{i}^{m}, \xi_{i}), \varphi(\xi_{i+1}^{m}, \xi_{i+1})\right) \\ = \frac{m}{2}\left(\varphi(\xi_{i}^{m}, \xi_{i}) + \varphi(\xi_{i+1}^{m}, \xi_{i+1})\right)^{\mathrm{T}}\left(\varphi(\xi_{i}^{m}, \xi_{i}) + \varphi(\xi_{i+1}^{m}, \xi_{i+1})\right), \quad (17)$$

where m > 0 is a scalar, and $\varphi(\xi_i^m, \xi_i)$ and $\varphi(\xi_{i+1}^m, \xi_{i+1})$ are according to (10). Note that the artificial kinetic energy T_i can also be written as the summation of the artificial rotational kinetic energy \mathcal{T}_i^r and artificial translational kinetic energy \mathcal{T}_i^r

employing (10) as

$$\mathcal{T}_{i} = \mathcal{T}^{r}(\omega_{i+1}, \omega_{i}) + \mathcal{T}^{t}(v_{i+1}, v_{i}) = \mathcal{T}_{i}^{r} + \mathcal{T}_{i}^{t}$$

= $\frac{1}{2}(\omega_{i+1} + \omega_{i})^{\mathrm{T}}m(\omega_{i+1} + \omega_{i}) + \frac{1}{2}(v_{i+1} + v_{i})^{\mathrm{T}}m(v_{i+1} + v_{i}).$ (18)

Let the discrete-time Lagrangian be defined as the difference between the artificial kinetic energy and artificial potential energy terms:

$$\begin{aligned} \mathscr{L}_{i} &:= \mathscr{L}(\omega_{i+1}, \omega_{i}, v_{i+1}, v_{i}, \hat{g}_{i}, L_{i}^{m}, D_{i}, \bar{a}_{i}^{m}, \bar{p}_{i}) \\ &= \mathcal{T}^{r}(\omega_{i+1}, \omega_{i}) + \mathcal{T}^{t}(v_{i+1}, v_{i}) - \mathcal{U}^{r}(\hat{g}_{i}, L_{i}^{m}, D_{i}) - \mathcal{U}^{t}(\hat{g}_{i}, \bar{a}_{i}^{m}, \bar{p}_{i}) \\ &= \frac{1}{2}(\omega_{i+1} + \omega_{i})^{\mathrm{T}}m(\omega_{i+1} + \omega_{i}) + \frac{1}{2}(v_{i+1} + v_{i})^{\mathrm{T}}m(v_{i+1} + v_{i}) \\ &- \frac{1}{2}k_{p}\langle D_{i} - \hat{R}_{i}L_{i}^{m}, (D_{i} - \hat{R}_{i}L_{i}^{m})W_{i}\rangle - \kappa ||\bar{p}_{i} - \hat{R}_{i}\bar{a}_{i}^{m} - \hat{b}_{i}||^{2}. \end{aligned}$$
(19)

If the estimation process is started at time t_0 , then the discrete-time action functional corresponding to the discrete-time Lagrangian (19) over the time interval $[t_0, T]$ can be expressed as

$$\mathfrak{s}_{d}(\mathscr{L}_{i}) := h \sum_{i=0}^{N} \left\{ \frac{1}{2} (\omega_{i+1} + \omega_{i})^{\mathrm{T}} m(\omega_{i+1} + \omega_{i}) + \frac{1}{2} (v_{i+1} + v_{i})^{\mathrm{T}} m(v_{i+1} + v_{i}) - \frac{1}{2} k_{p} \langle D_{i} - \hat{R}_{i} L_{i}^{m}, (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \rangle - \kappa ||\bar{p}_{i} - \hat{R}_{i} \bar{a}_{i}^{m} - \hat{b}_{i}||^{2} \right\}.$$
(20)

We are now ready to state our first result on optimal pose estimation.

Proposition 1 The variational filter for pose minimizing the action functional $\mathfrak{s}_d(\mathscr{L}_i)$ defined in (20) is given as

$$\begin{cases} \varphi_{i+2} + \varphi_{i+1} = \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \left[(\varphi_{i+1} + \varphi_{i}) - \frac{h}{2m}Z'_{i} - \frac{h}{2m}\eta_{i+1}\right] \\ \hat{\xi}_{i} = \xi_{i}^{m} - \varphi_{i}, \\ \hat{g}_{i+1} = \hat{g}_{i} \begin{bmatrix} \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \frac{h}{2}(\hat{\nu}_{i} + \hat{\nu}_{i+1}) \\ 0 & 1 \end{bmatrix}, \end{cases}$$
(21)

where η_{i+1} contains Rayleigh dissipation terms for angular and translational motions defined as

$$\eta_{i+1} := \begin{bmatrix} \tau_{i+1} \\ f_{i+1} \end{bmatrix},$$

with $Z'_i := Z'(\hat{g}_{i+1}, \hat{g}_i, L^m_{i+1}, D_{i+1}, \bar{a}^m_{i+1}, \tau_i, f_i)$ defined by

$$Z'_{i} := \begin{bmatrix} -k_{p} S_{\Gamma_{i+1}}(\hat{R}_{i+1}) + m(\hat{v}_{i+1} + v_{i})^{\times}(v_{i+1} + v_{i}) + \kappa(\bar{a}^{m}_{i+1})^{\times}\hat{R}^{\mathrm{T}}_{i+1}y_{i+1} \\ \kappa \hat{R}^{\mathrm{T}}_{i+1}y_{i+1}, \end{bmatrix}$$

🖄 Springer

where $S_{\Gamma_i}(\hat{R}_i) := vex(\Gamma_i^{\mathrm{T}}\hat{R}_i - \hat{R}_i^{\mathrm{T}}\Gamma_i)$ and $\Gamma_i := D_i W_i(L_i^m)^{\mathrm{T}}$.

Proof Consider a first variation in the discrete attitude estimate as

$$\delta \hat{R}_i = \hat{R}_i \Sigma_i^{\times}, \tag{22}$$

where $\Sigma_i \in \mathbb{R}^3$ represents a variation for the discrete attitude estimate. For fixed endpoint variations, we have $\Sigma_0 = \Sigma_N = 0$. A first-order approximation is to assume that $\hat{\Omega}^{\times}$ and $\delta \hat{\Omega}^{\times}$ commute. Taking the first variation of the discrete-time attitude kinematics according to the first equation of (6) and comparing with (22), we get

$$\delta \hat{R}_{i+1} = \delta \hat{R}_i \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) + \frac{h}{2}\hat{R}_i \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) \delta(\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times} = \hat{R}_{i+1}\Sigma_{i+1}^{\times}.$$
(23)

The above can be rearranged to

$$\hat{R}_{i+1}\frac{h}{2}\delta(\hat{\Omega}_{i+1}+\hat{\Omega}_{i})^{\times} = \hat{R}_{i+1}\Sigma_{i+1}^{\times}$$
$$-\hat{R}_{i+1}\operatorname{Ad}_{\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1}+\hat{\Omega}_{i})^{\times}\right)}\Sigma_{i}^{\times}$$
$$\Rightarrow \frac{h}{2}\delta(\hat{\Omega}_{i+1}+\hat{\Omega}_{i})^{\times} = \Sigma_{i+1}^{\times} - \operatorname{Ad}_{\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1}+\hat{\Omega}_{i})^{\times}\right)}\Sigma_{i}^{\times}, \qquad (24)$$

which can be equivalently written as an equation in \mathbb{R}^3 as follows:

$$\frac{h}{2}\delta(\hat{\Omega}_{i+1}+\hat{\Omega}_i) = \Sigma_{i+1} - \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1}+\hat{\Omega}_i)^{\times}\right)\Sigma_i.$$
(25)

Taking the variation of $\omega_i = \Omega_i^m - \hat{\Omega}_i$

$$\delta(\omega_{i+1} + \omega_i) = -\delta(\hat{\Omega}_{i+1} + \hat{\Omega}_i).$$
⁽²⁶⁾

Next, we assume the variation in \hat{b}_i to be,

$$\delta \hat{b}_i = \hat{R}_i \rho_i, \tag{27}$$

where $\rho_i \in \mathbb{R}^3$ represents the variation in the discrete position estimate. For fixed endpoint variations, we have $\rho_0 = \rho_N = 0$. Taking the first variation of the discrete-time position kinematics according to the second equation of (6) and comparing with (27), we get

$$\begin{split} \delta \hat{b}_{i+1} &= \delta \hat{b}_i + \delta \hat{R}_{i+1} \frac{h}{2} (\hat{v}_{i+1} + \hat{v}_i) + \hat{R}_{i+1} \frac{h}{2} \delta (\hat{v}_{i+1} + \hat{v}_i) \\ &\Rightarrow \hat{R}_{i+1} \rho_{i+1} = \hat{R}_i \rho_i + \hat{R}_{i+1} \sum_{i+1}^{\times} \frac{h}{2} (\hat{v}_{i+1} + \hat{v}_i) + \hat{R}_{i+1} \frac{h}{2} \delta (\hat{v}_{i+1} + \hat{v}_i) \\ &\Rightarrow \frac{h}{2} \delta (\hat{v}_{i+1} + \hat{v}_i) = \rho_{i+1} - \exp\left(-\frac{h}{2} (\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) \rho_i - \sum_{i+1}^{\times} \frac{h}{2} (\hat{v}_{i+1} + \hat{v}_i) \end{split}$$

$$(28)$$

and the variation of $v_i = v_i^m - \hat{v}_i$ gives us

$$\delta(v_{i+1} + v_i) = -\delta(\hat{v}_{i+1} + \hat{v}_i).$$
⁽²⁹⁾

We have $y_i = \bar{p}_i - \hat{R}_i \bar{a}_i^m - \hat{b}_i$. Therefore,

$$\delta y_i = -\delta \hat{R}_i \bar{a}_i^m - \delta \hat{b}_i$$

= $\hat{R}_i \Sigma_i^{\times} \bar{a}_i^m - \hat{R}_i \rho_i$
= $\hat{R}_i \left(\left(\bar{a}_i^m \right)^{\times} \Sigma_i - \rho_i \right).$ (30)

Consider the artificial potential energy term in (14). Taking its first variation with respect to the estimated attitude \hat{R} , we get

$$\begin{split} \delta \mathcal{U}_{i}^{r} &= \frac{k_{p}}{2} \left\{ \langle -\delta \hat{R}_{i} L_{i}^{m}, (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \rangle + \langle D_{i} - \hat{R}_{i} L_{i}^{m}, (-\delta \hat{R}_{i} L_{i}^{m}) W_{i} \rangle \right\} \\ &= k_{p} \langle -\delta \hat{R}_{i} L_{i}^{m}, (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \rangle \\ &= k_{p} \langle -\hat{R}_{i} \Sigma_{i}^{\times}, (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \rangle \\ &= k_{p} \operatorname{tr} \left((L_{i}^{m})^{\mathrm{T}} \Sigma_{i}^{\times} \hat{R}_{i}^{\mathrm{T}} (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \right) \\ &= k_{p} \operatorname{tr} \left((\Sigma_{i}^{\times})^{\mathrm{T}} L_{i}^{m} W_{i} D_{i}^{\mathrm{T}} \hat{R}_{i} \right) \\ &= k_{p} \left\{ \Sigma_{i}^{\times}, L_{i}^{m} W_{i} D_{i}^{\mathrm{T}} \hat{R}_{i} - \hat{R}_{i}^{\mathrm{T}} D_{i} W_{i} (L_{i}^{m})^{\mathrm{T}} \right\} \\ &= k_{p} \frac{1}{2} \langle \Sigma_{i}^{\times}, \Gamma_{i}^{\mathrm{T}} \hat{R}_{i} - \hat{R}_{i}^{\mathrm{T}} \Gamma_{i} \rangle = k_{p} S_{\Gamma_{i}}^{\mathrm{T}} (\hat{R}_{i}) \Sigma_{i}. \end{split}$$

Similarly, taking the first variation of the artificial potential energy function in (15) and using results from (30), we get

$$\delta \mathcal{U}_i^t = \kappa y_i^{\mathrm{T}} \delta y_i = \kappa y_i^{\mathrm{T}} \left(\left(\bar{a}_i^m \right)^{\times} \Sigma_i - \rho_i \right).$$
(32)

Similarly, we also obtain the variations in artificial kinetic energies as follows

$$\delta \mathcal{T}_i^r = \frac{2m}{h} (\omega_{i+1} + \omega_i)^{\mathrm{T}} \left(\exp\left(-\frac{h}{2} (\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) \Sigma_i - \Sigma_{i+1} \right), \tag{33}$$

$$\delta \mathcal{T}_{i}^{t} = \frac{2m}{h} (v_{i+1} + v_{i})^{\mathrm{T}} \left(\exp\left(-\frac{h}{2} (\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \rho_{i} + \Sigma_{i+1}^{\times} \frac{h}{2} (\hat{v}_{i+1} + \hat{v}_{i}) - \rho_{i+1} \right),$$
(34)

with the help of relations from (25) to (29). Taking the first variation of the discretetime action sum in (20) and employing (31) to (34), we obtain

$$\delta \mathfrak{s}_{d} = h \sum_{i=0}^{N} \left\{ \delta \mathcal{T}_{i}^{r} + \delta \mathcal{T}_{i}^{t} - \delta \mathcal{U}_{i}^{r} - \delta \mathcal{U}_{i}^{t} \right\}$$

$$= h \sum_{i=0}^{N} \left\{ \frac{2m}{h} (\omega_{i+1} + \omega_{i})^{\mathrm{T}} \left(\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \Sigma_{i} - \Sigma_{i+1} \right) + \frac{2m}{h} (\upsilon_{i+1} + \upsilon_{i})^{\mathrm{T}} \left(\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \rho_{i} + \Sigma_{i+1}^{\times} \frac{h}{2}(\hat{\upsilon}_{i+1} + \hat{\upsilon}_{i}) - \rho_{i+1} \right) - k_{p} S_{\Gamma_{i}}^{\mathrm{T}}(\hat{R}_{i}) \Sigma_{i} - \kappa y_{i}^{\mathrm{T}} \left(\left(\bar{a}_{i}^{m} \right)^{\times} \Sigma_{i} - \rho_{i} \right) \right\}.$$
(35)

We now apply the discrete Lagrange–d'Alembert principle (Marsden and West 2001) with two Rayleigh dissipation terms $\tau_i \in \mathbb{R}^3$ and $f_i \in \mathbb{R}^3$ for angular and translational motion, respectively,

$$\delta \mathfrak{s}_{d} + h \sum_{i=0}^{N-1} \left\{ \tau_{i}^{\mathrm{T}} \Sigma_{i} + f_{i}^{\mathrm{T}} \rho_{i} \right\} = 0$$

$$\Rightarrow \sum_{i=0}^{N-1} \left\{ 2m(\omega_{i+1} + \omega_{i})^{\mathrm{T}} \left(\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \Sigma_{i} - \Sigma_{i+1}\right) + 2m(\upsilon_{i+1} + \upsilon_{i})^{\mathrm{T}} \left(\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_{i})^{\times}\right) \rho_{i} + \Sigma_{i+1}^{\times} \frac{h}{2}(\hat{\upsilon}_{i+1} + \hat{\upsilon}_{i}) - \rho_{i+1}\right) - k_{p}hS_{\Gamma_{i}}^{\mathrm{T}}(\hat{R}_{i})\Sigma_{i} - \kappa hy_{i}^{\mathrm{T}} \left(\left(\bar{a}_{i}^{m}\right)^{\times} \Sigma_{i} - \rho_{i}\right) h\tau_{i}^{\mathrm{T}}\Sigma_{i} + hf_{i}^{\mathrm{T}}\rho_{i} \right\} = 0.$$
(36)

Splitting (36) into two equations assuming independence of Σ_i and ρ_i will give us

$$2m(\omega_{i+2} + \omega_{i+1})^{\mathrm{T}} \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) - 2m(\omega_{i+1} + \omega_{i})^{\mathrm{T}} - hm(v_{i+1} + v_{i})^{\mathrm{T}}(\hat{v}_{i+1} + \hat{v}_{i})^{\times} - -k_{p}hS_{\Gamma_{i+1}}^{\mathrm{T}}(\hat{R}_{i+1}) - \kappa hy_{i+1}^{\mathrm{T}}\left(\bar{a}_{i+1}^{m}\right)^{\times} + h\tau_{i+1}^{\mathrm{T}} = 0,$$
(37)

and

$$2m(v_{i+2} + v_{i+1})^{\mathrm{T}} \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) - 2m(v_{i+1} + v_i)^{\mathrm{T}} + \kappa h y_{i+1}^{\mathrm{T}} \hat{R}_{i+1} + h f_i^{\mathrm{T}} = 0.$$
(38)

The above equations can be simplified to obtain

$$\omega_{i+2} + \omega_{i+1} = \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \left[\omega_{i+1} + \omega_i + \frac{h}{2m} \{k_p S_{\Gamma_{i+1}}(\hat{R}_{i+1}) - m(\hat{v}_{i+1} + v_i)^{\times}(v_{i+1} + v_i) - \kappa(\bar{a}_{i+1}^m)^{\times}\hat{R}_{i+1}^{\mathsf{T}}y_{i+1} - \tau_{i+1}\}\right] = 0,$$
(39)

and

$$v_{i+2} + v_{i+1} = \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \left[v_{i+1} + v_i - \frac{h}{2m}\left\{\kappa \hat{R}_{i+1}^{\mathrm{T}} y_{i+1} + f_{i+1}\right\}\right] = 0,$$
(40)

which in turn can be combined to obtain (21).

5 Discrete-Time Asymptotically Stable and Optimal Pose Estimator

For the discrete-time Lyapunov analysis, we use the same combination of artificial potential energy-like terms as defined in (16). We construct a new kinetic energy-like term to encapsulate the error in the generalized velocity estimation. This term will aid us in the Lyapunov analysis. We propose the following kinetic energy-like term:

$$\mathcal{T}_i^l := \mathcal{T}^l \left(\varphi(\xi_i^m, \xi_i) \right) = \frac{m}{2} \varphi(\xi_i^m, \xi_i)^{\mathrm{T}} \varphi(\xi_i^m, \xi_i), \tag{41}$$

where m > 0 is a scalar. We carry out the Lyapunov analysis in the absence of measurement errors. The following Lemma provides the form of the artificial potential energy term in the absence of measurement errors.

Lemma 1 In the absence of measurement noise, the artificial potential energy defined in (16) takes the form

$$\mathcal{U}_i = \mathcal{U}(\hat{g}_i, L_i^m, D_i, \bar{a}_i^m, \bar{p}_i) = k_p \langle I - Q_i, K_i \rangle + \kappa y_i^{\mathrm{T}} y_i, \qquad (42)$$

where $K_i = D_i W_i D_i^{\mathrm{T}}$ and $y_i = y(h_i, \bar{p}_i) = Q_i^{\mathrm{T}} x_i + (I - Q_i^{\mathrm{T}}) \bar{p}_i$.

Proof In the absence of measurements errors, we have $L_i^m = L_i$, $\bar{a}_i^m = \bar{a}_i$ and $\xi_i^m = \xi_i$. The rotational potential function can be rewritten as

$$\mathcal{U}^{r}(\hat{g}_{i}, L_{i}^{m}, D_{i}) = \frac{k_{p}}{2} \langle D_{i} - \hat{R}_{i} L_{i}^{m}, (D_{i} - \hat{R}_{i} L_{i}^{m}) W_{i} \rangle$$

$$= \frac{k_{p}}{2} \langle I - R_{i} \hat{R}_{i}^{T}, D_{i} W_{i} D_{i}^{T} \rangle$$

$$\Rightarrow \mathcal{U}^{r}(Q_{i}) = k_{p} \langle I - Q_{i}, K_{i} \rangle \text{ where } K_{i} = D_{i} W_{i} D_{i}^{T}.$$

🖄 Springer

Since we have $\hat{R}_i L_i = Q_i^T D_i$ and $\hat{b}_i = Q_i^T (b_i - x_i)$ we get

$$y_i = y(h_i, \bar{p}_i) = \bar{p}_i - \hat{R}_i \bar{a}_i - \hat{b}_i = Q_i^{\mathrm{T}} x_i + (I - Q_i^{\mathrm{T}}) \bar{p}_i,$$

as the form of potential energy in the absence of measurement errors.

We state the following assumption that is relevant for the proof of asymptotic stability.

Assumption 1 The measured beacons and inertial vectors are fixed throughout the estimation process which results in $K_i = K$ and $\bar{p}_i = \bar{p}$ for some constants $K \in \mathbb{R}^{n \times n}$ and $\bar{p} \in \mathbb{R}^3$ for all t_i .

We are now ready to state the main result of this article on an optimal asymptotically stable pose filter.

Theorem 2 Consider the following form of the dissipation term in (43)

$$\eta_{i+1} = \frac{2}{h} \left\{ m(\varphi_{i+1} + \varphi_i) - \frac{h}{2} Z'_i - \frac{m}{m+l} \exp\left(\frac{h}{2} (\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \left[2m\varphi_{i+1} - hZ_{i+1}\right] \right\}.$$
 (43)

Then the resulting nonlinear pose estimator given by,

$$\begin{cases} \varphi_{i+1} = \frac{1}{m+l} \left[(m-l)\varphi_i - hZ_i \right] \\ \hat{\xi}_i = \xi_i^m - \varphi_i, \\ \hat{g}_{i+1} = \hat{g}_i \begin{bmatrix} \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) \exp\left(\frac{h}{2}(\hat{\Omega}_{i+1} + \hat{\Omega}_i)^{\times}\right) \frac{h}{2} \left(\hat{\nu}_i + \hat{\nu}_{i+1}\right) \\ 0 & 1 \end{bmatrix}, \end{cases}$$
(44)

where and $Z_i = Z(\hat{g}_i, L_i^m, D_i, \bar{a}_i^m, \bar{p}_i)$ is defined by

$$Z(\hat{g}_{i}, L_{i}^{m}, D_{i}, \bar{a}_{i}^{m}, \bar{p}_{i}) = \begin{bmatrix} -k_{p}S_{\Gamma_{i}}(\hat{R}_{i}) + \kappa \hat{R}_{i}^{\mathrm{T}} \left(Q_{i}^{\mathrm{T}}(\bar{p} - b_{i})\right)^{\times} (y_{i+1} + y_{i}) \\ \kappa \hat{R}_{i+1}^{\mathrm{T}}(y_{i+1} + y_{i}) \end{bmatrix}$$
(45)

where $\Gamma_i = D_i W_i (L_i^m)^T$, $S_{\Gamma_i}(\hat{R}_i) = vex(\Gamma_i^T \hat{R}_i - \hat{R}_i^T \Gamma_i)$ and $l > 0, l \neq m$, is asymptotically stable under Assumption 1 at the estimation error state $(\bar{g}_i, \varphi_i) = (I, 0)$. Further, the domain of attraction of $(\bar{g}_i, \varphi_i) = (I, 0)$ is a dense open subset of SE(3) × \mathbb{R}^6 .

Proof We choose the following discrete-time Lyapunov candidate:

$$V_i := V(Q_i, h_i, \bar{p}_i) := \mathcal{U}_i + \mathcal{T}_i,$$

The stability of the attitude and angular velocity error can be shown by analyzing $\Delta V_i = \Delta U_i + \Delta T_i$.

Using Assumption 1, we first calculate

$$\Delta \mathcal{U}_{i}^{r} = \mathcal{U}_{i+1}^{r} - \mathcal{U}_{i}^{r} = k_{p} \langle I - Q_{i+1}, K \rangle - k_{p} \langle I - Q_{i}, K \rangle,$$

$$\Delta \mathcal{U}_{i}^{r} = k_{p} \langle Q_{i} - Q_{i+1}, K \rangle = -k_{p} \langle \Delta Q_{i}, K \rangle,$$

where $\Delta Q_i = Q_{i+1} - Q_i$. Now,

$$\Delta Q_i = Q_{i+1} - Q_i$$

= $Q_i \left[\hat{R}_i \exp\left(\frac{h}{2}(\hat{\omega}_{i+1} + \hat{\omega}_i)^{\times}\right) \hat{R}_i^{\mathrm{T}} - I \right]$

Approximating $\exp\left(\frac{h}{2}(\hat{\omega}_{i+1}+\hat{\omega}_i)^{\times}\right)$ as shown in (12), we have

$$\Delta Q_i = Q_i \left[\hat{R}_i \left(I + \frac{h}{2} (\hat{\omega}_{i+1} + \hat{\omega}_i)^{\times} \right) \hat{R}_i^{\mathrm{T}} - I \right]$$

$$= \frac{h}{2} Q_i \left(\hat{R}_i (\hat{\omega}_{i+1} + \hat{\omega}_i)^{\times} \hat{R}_i^{\mathrm{T}} \right)$$

$$= \frac{h}{2} Q_i \left(\hat{R}_i (\hat{\omega}_{i+1} + \hat{\omega}_i) \right)^{\times}.$$

In the absence of measurement errors, we have $L_i^m = R_i^T D_i$. Therefore,

$$\begin{split} \Delta \mathcal{U}_i^r &= -\frac{k_p h}{2} \left\langle Q_i \left(\hat{R}_i \left(\omega_{i+1} + \omega_i \right) \right)^{\times}, K_i \right\rangle \\ &= -\frac{k_p h}{2} \left\langle R_i (\omega_{i+1} + \omega_i)^{\times} \hat{R}_i^{\mathrm{T}}, D_i W_i D_i^{\mathrm{T}} \right\rangle \\ &= -\frac{k_p h}{2} \left\langle (\omega_{i+1} + \omega_i)^{\times} \hat{R}_i^{\mathrm{T}}, R_i^{\mathrm{T}} D_i W_i D_i^{\mathrm{T}} \right\rangle \\ &= -\frac{k_p h}{2} \left\langle (\omega_{i+1} + \omega_i)^{\times} \hat{R}_i^{\mathrm{T}}, L_i^m W_i D_i^{\mathrm{T}} \right\rangle. \end{split}$$

We can further simplify the above expression using $\Gamma_i = D_i W_i (L_i^m)^T$.

$$\Delta \mathcal{U}_{i}^{r} = -\frac{k_{p}h}{2} \left\langle (\omega_{i+1} + \omega_{i})^{\times}, \Gamma_{i}^{\mathrm{T}}\hat{R}_{i} \right\rangle$$
$$= -\frac{k_{p}h}{4} \left\langle (\omega_{i+1} + \omega_{i})^{\times}, \Gamma_{i}^{\mathrm{T}}\hat{R}_{i} - \hat{R}_{i}^{\mathrm{T}}\Gamma_{i} \right\rangle$$
$$= -\frac{k_{p}h}{2} (\omega_{i+1} + \omega_{i})^{\mathrm{T}}S_{\Gamma_{i}}(\hat{R}_{i})$$
$$= -\frac{h}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \begin{bmatrix} k_{p}S_{\Gamma_{i}}(\hat{R}_{i}) \\ 0 \end{bmatrix}$$
(46)

where $S_{\Gamma_i}(\hat{R}_i) = \operatorname{vex}(\Gamma_i^{\mathrm{T}}\hat{R}_i - \hat{R}_i^{\mathrm{T}}\Gamma_i).$

D Springer

Similarly we can compute the change in the translational potential energy as follows:

$$\Delta \mathcal{U}_i^t = \mathcal{U}^t(y_{i+1}) - \mathcal{U}^t(y_i)$$

= $(y_{i+1} - y_i)^{\mathrm{T}} \kappa(y_{i+1} + y_i)$
$$\Delta \mathcal{U}_i^t = (y_{i+1} - y_i)^{\mathrm{T}} \kappa(y_{i+1} + y_i).$$

Under Assumption 1, we have,

$$y_{i} = Q_{i}^{\mathrm{T}} x_{i} + (I - Q_{i}^{\mathrm{T}})\bar{p}$$

$$y_{i+1} = Q_{i+1}^{\mathrm{T}} x_{i+1} + (I - Q_{i+1}^{\mathrm{T}})\bar{p}.$$
(47)

Therefore, the discrete-time evolution of y_i is

$$y_{i+1} - y_i = Q_{i+1}^{\mathrm{T}} b_{i+1} - \hat{b}_{i+1} - \left(Q_i^{\mathrm{T}} b_i - \hat{b}_i\right) + \left(Q_i - Q_{i+1}\right)^{\mathrm{T}} \bar{p}$$

$$= \hat{R}_{i+1} R_{i+1}^{\mathrm{T}} \left(b_i + R_{i+1} \frac{h}{2} \left(v_i + v_{i+1}\right)\right) - Q_i^{\mathrm{T}} b_i$$

$$- \left(\hat{b}_i + \hat{R}_{i+1} \frac{h}{2} \left(\hat{v}_i + \hat{v}_{i+1}\right) - \hat{b}_i\right)$$

$$= \left(Q_{i+1} - Q_i\right)^{\mathrm{T}} \left(b_i - \bar{p}\right) + \hat{R}_{i+1} \frac{h}{2} \left(v_i + v_{i+1}\right)$$

$$= \frac{h}{2} \left[Q_i \left(\hat{R}_i (\omega_{i+1} + \omega_i)\right)^{\times}\right]^{\mathrm{T}} \left(b_i - \bar{p}\right) + \hat{R}_{i+1} \frac{h}{2} \left(v_i + v_{i+1}\right)$$

$$= -\frac{h}{2} \left(\hat{R}_i (\omega_{i+1} + \omega_i)\right)^{\times} Q_i^{\mathrm{T}} \left(b_i - \bar{p}\right) + \hat{R}_{i+1} \frac{h}{2} \left(v_i + v_{i+1}\right)$$

$$= \frac{h}{2} \left(Q_i^{\mathrm{T}} \left(b_i - \bar{p}\right)\right)^{\times} \hat{R}_i \left(\omega_{i+1} + \omega_i\right) + \hat{R}_{i+1} \frac{h}{2} \left(v_i + v_{i+1}\right). \quad (48)$$

Using (48), we obtain

$$\Delta \mathcal{U}_{i}^{t} = (y_{i+1} - y_{i})^{\mathrm{T}} \kappa (y_{i+1} + y_{i})
= \frac{\kappa h}{2} \left[(\omega_{i+1} + \omega_{i})^{\mathrm{T}} \hat{R}_{i}^{\mathrm{T}} \left(\mathcal{Q}_{i}^{\mathrm{T}} (\bar{p} - b_{i}) \right)^{\times} + (v_{i} + v_{i+1})^{\mathrm{T}} \hat{R}_{i+1}^{\mathrm{T}} \right] (y_{i+1} + y_{i})
= \frac{h}{2} \left[\frac{\omega_{i+1} + \omega_{i}}{v_{i} + v_{i+1}} \right]^{\mathrm{T}} \left[\kappa \hat{R}_{i}^{\mathrm{T}} \left(\mathcal{Q}_{i}^{\mathrm{T}} (\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ \kappa \hat{R}_{i+1}^{\mathrm{T}} (y_{i+1} + y_{i}) \right]
= \frac{h}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \left[\kappa \hat{R}_{i}^{\mathrm{T}} \left(\mathcal{Q}_{i}^{\mathrm{T}} (\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ \kappa \hat{R}_{i+1}^{\mathrm{T}} (y_{i+1} + y_{i}) \right].$$
(49)

Similarly we can compute the change in the kinetic energy as follows:

$$\Delta T_i = T(\varphi_{i+1}) - T(\varphi_i)$$

$$= (\varphi_{i+1} + \varphi_i)^{\mathrm{T}} \frac{m}{2} (\varphi_{i+1} - \varphi_i).$$
(50)

Using values from (46), (49) and (50), we obtain

$$\begin{aligned} \Delta V_{i} &= \Delta U_{i}^{r} + \Delta U_{i}^{i} + \Delta T_{i} \\ &= -\frac{h}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \begin{bmatrix} k_{p} S_{\Gamma_{i}}(\hat{R}_{i}) \\ 0 \end{bmatrix} + (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \begin{bmatrix} \kappa \hat{R}_{i}^{\mathrm{T}} \left(Q_{i}^{\mathrm{T}}(\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ &+ \frac{h}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \begin{bmatrix} \kappa \hat{R}_{i}^{\mathrm{T}} \left(Q_{i}^{\mathrm{T}}(\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ &- \kappa \hat{R}_{i+1}^{\mathrm{T}}(y_{i+1} + y_{i}) \end{bmatrix} \end{aligned}$$

$$= \frac{1}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \left\{ m(\varphi_{i+1} - \varphi_{i}) + \kappa \hat{R}_{i+1}^{\mathrm{T}} \left(Q_{i}^{\mathrm{T}}(\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ &+ h \begin{bmatrix} -k_{p} S_{\Gamma_{i}}(\hat{R}_{i}) + \kappa \hat{R}_{i+1}^{\mathrm{T}} \left(Q_{i}^{\mathrm{T}}(\bar{p} - b_{i}) \right)^{\times} (y_{i+1} + y_{i}) \\ &- \kappa \hat{R}_{i+1}^{\mathrm{T}}(y_{i+1} + y_{i}) \end{bmatrix} \right\}. \tag{51}$$

Using (45) yields

$$\Delta V_{i} = \frac{1}{2} (\varphi_{i+1} + \varphi_{i})^{\mathrm{T}} \{ m(\varphi_{i+1} - \varphi_{i}) + hZ_{i} \}.$$

Therefore, ΔV_{i+1} can be written as

$$\Delta V_{i+1} = \frac{1}{2} (\varphi_{i+2} + \varphi_{i+1})^{\mathrm{T}} \{ m(\varphi_{i+2} - \varphi_{i+1}) + hZ_{i+1} \}.$$

Substituting for φ_{i+2} from (21),

$$\Delta V_{i+1} = \frac{1}{2} (\varphi_{i+2} + \varphi_{i+1})^{\mathrm{T}} \bigg\{ -2m\varphi_{i+1} + hZ_{i+1} + \exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \bigg[m(\varphi_{i+1} + \varphi_i) - \frac{h}{2}Z'_i - \frac{h}{2}\eta_{i+1} \bigg] \bigg\}.$$

Now, in order for ΔV_{i+1} to be negative definite, we require

$$\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2}+\hat{\Omega}_{i+1})^{\times}\right)\left[m(\varphi_{i+1}+\varphi_{i})-\frac{h}{2}Z'_{i}-\frac{h}{2}\eta_{i+1}\right] -2m\varphi_{i+1}+hZ_{i+1}=-l(\varphi_{i+2}+\varphi_{i+1}),$$
(52)

where $l > 0, l \neq m$. ΔV_{i+1} simplifies to

$$\Delta V_{i+1} = -\frac{l}{2} \left(\varphi_{i+2} + \varphi_{i+1} \right)^{\mathrm{T}} \left(\varphi_{i+2} + \varphi_{i+1} \right).$$
(53)

D Springer

Again substituting the value of φ_{i+2} from (21) into (52), we have

$$\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2}+\hat{\Omega}_{i+1})^{\times}\right)\left[m(\varphi_{i+1}+\varphi_{i})-\frac{h}{2}Z'_{i}-\frac{h}{2}\eta_{i+1}\right]-2m\varphi_{i+1} + hZ_{i+1} = -\frac{l}{m}\exp\left(-\frac{h}{2}(\hat{\Omega}_{i+2}+\hat{\Omega}_{i+1})^{\times}\right)\left[m(\varphi_{i+1}+\varphi_{i})-\frac{h}{2}Z'_{i}-\frac{h}{2}\eta_{i+1}\right].$$

Streamlining the terms in the above equation, we obtain

$$\frac{m+l}{m} \exp \left(\frac{h}{2} (\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times} \right) \left[m(\varphi_{i+1} + \varphi_i) - \frac{h}{2} Z'_i - \frac{h}{2} \eta_{i+1} \right]$$

= $2m\varphi_{i+1} - hZ_{i+1},$

which upon further simplification yields

$$m(\varphi_{i+1} + \varphi_i) - \frac{h}{2} Z'_i - \frac{h}{2} \eta_{i+1} = \frac{m}{m+l} \exp\left(\frac{h}{2} (\hat{\Omega}_{i+2} + \hat{\Omega}_{i+1})^{\times}\right) \left[2m\varphi_{i+1} - hZ_{i+1}\right].$$

Rearranging the terms above, we obtain the value of η_i as shown in (43). After substituting for η_i in (44), we get

$$\varphi_{i+2} = \frac{1}{m+l} \left[(m-l)\varphi_{i+2} - hZ_{i+1} \right].$$
(54)

(54) can be rewritten in the previous time step as

$$\varphi_{i+1} = \frac{1}{m+l} \left[(m-l)\varphi_i - hZ_i \right].$$
(55)

We can rewrite ΔV_i with the help of (53) to be,

$$\Delta V_i = -\frac{l}{2} \left(\varphi_{i+1} + \varphi_i\right)^{\mathrm{T}} \left(\varphi_{i+1} + \varphi_i\right).$$
(56)

We employ the discrete-time La Salle invariance principle from LaSalle (1976) considering our domain (SE(3) $\times \mathbb{R}^6$) to be a subset of \mathbb{R}^{12} . We use Theorem 6.3 and Theorem 7.9 from Chapter 1 of LaSalle (1976). For this, we first compute,

$$\mathscr{E} = \Delta V_i^{-1}(0) = \{ (\bar{g}_i, \varphi_i) \in SE(3) \times \mathbb{R}^6 \mid \varphi_{i+1} + \varphi_i \equiv 0 \}.$$
(57)

Now,

$$\varphi_{i+1} + \varphi_i = 0 \Rightarrow \omega_{i+1} + \omega_i = 0, \quad v_{i+1} + v_i = 0.$$
 (58)

From (13), we obtain that $x_{i+1} = x_i$, and therefore, from (47) we also have that $y_{i+1} = y_i$. Note that we have $\bar{g}_{i+1} = \bar{g}_i$ when $Q_{i+1} = Q_i$ and $x_{i+1} = x_i$. Further substituting (58) into (44), we have

$$\varphi_{i} = \frac{hZ_{i}}{2m} = \frac{h\kappa}{m} \begin{bmatrix} \hat{R}_{i}^{\mathrm{T}} \left(\mathcal{Q}_{i}^{\mathrm{T}} (\bar{p} - b_{i}) \right)^{\times} (y_{i}) \\ \hat{R}_{i+1}^{\mathrm{T}} (y_{i}) \end{bmatrix}$$

Similarly,

$$\varphi_{i+1} = \frac{h\kappa}{m} \begin{bmatrix} \hat{R}_{i+1}^{\mathrm{T}} \left(\mathcal{Q}_{i}^{\mathrm{T}} (\bar{p} - b_{i+1}) \right)^{\times} (y_{i}) \\ \hat{R}_{i+2}^{\mathrm{T}} (y_{i}) \end{bmatrix}$$

and since $\varphi_{i+1} + \varphi_i = 0$, we have $(\hat{R}_{i+1} + \hat{R}_{i+2})^T y_i = 0$ which can be rewritten as $(I + \hat{R}_{i+1}\hat{R}_{i+2}^T)y_i = 0$. It is evident that $(I + \hat{R}_{i+1}\hat{R}_{i+2}^T)$ is non-singular, and therefore, $y_i = 0$. Note that $y_i = Q_i^T x_i + (I - Q_i^T)\bar{p}$. Therefore, when $Q_i = I$, we have $y_i = x_i$ and subsequently $x_i = 0$ with $Q_i = I$ gives us $\hat{b}_i = b_i$. Therefore, the largest invariant set for the estimation error dynamics will be $\mathcal{M} = \{(\bar{g}_i, \varphi_i) \in$ SE(3) × $\mathbb{R}^6 | \bar{g}_i = I, \varphi_i = 0\}$. Furthermore, we obtain the positive limit set as the set,

$$\mathcal{I} := \mathcal{M} \cap V_i^{-1}(0)$$

= { $(\bar{g}_i, \varphi_i) \in \operatorname{SE}(3) \times \mathbb{R}^6 \mid \bar{g}_i = I, \varphi_i = 0$ }.

with an almost global domain of attraction.

This completes the proof of asymptotic stability of estimation error state $(\bar{g}_i, \varphi_i) = (I, 0)$ with an almost global domain of attraction.

6 Simulation Results

In order to numerically verify this estimator, simulated true states of an aerial vehicle are produced using a six-degree-of-freedom (6DOF) rigid body dynamics model. The continuous-time 6DOF rigid body dynamics equations are as follows:

$$\dot{b} = Rv$$

$$m_v \dot{v} = -m_v \Omega^{\times} v + \phi_v$$

$$\dot{R} = R\Omega^{\times}$$

$$J_v \dot{\Omega} = -\Omega^{\times} J_v \Omega + \tau_v,$$
(59)

where m_v and J_v are the mass and moment of inertia of the rigid body, respectively. ϕ_v and τ_v are the total force and torque in the body frame, respectively. A numerical method for the simulation of 6DOF rigid body dynamics is presented in Baraff (1997). Several efficient algorithms to compute rigid body dynamics can be found in Featherstone (2014). The vehicle mass and moment of inertia are taken to be $m_v = 0.42$ kg and $J_v = \text{diag}(10^{-3} \times [51.2 \ 60.2 \ 59.6])$ kg·m², respectively. The resultant external forces and torques applied on the vehicle are $\phi_v(t) = 10^{-3}[10\cos(0.1t)\ 2\sin(0.2t)\ -2\sin(0.5t)]^T$ N and $\tau_v(t) = 10^{-6}\phi_v(t)$ N·m, respectively. The flight area is assumed to be a cubic space of size $20 \times 20 \times 20$ with the origin of the inertial frame located at the center of this cube. The initial attitude and position of the vehicle are

$$R_0 = \exp_{\mathrm{SO}(3)} \left(\left(\frac{\pi}{4} \times \left[\frac{3}{7} - \frac{6}{7} \frac{2}{7} \right]^{\mathrm{T}} \right)^{\times} \right),$$

and $b_0 = [2.5 \ 0.5 \ -3]^{\mathrm{T}} \mathrm{m}.$

The vehicle's initial angular and translational velocities are

$$\Omega_0 = [0.2 - 0.05 \ 0.1]^{\mathrm{T}} \mathrm{rad/s},$$

and $\nu_0 = [-0.05 \ 0.15 \ 0.03]^{\mathrm{T}} \mathrm{m/s}$

The vehicle dynamics is simulated over a time interval of T = 60 s, with a time step size of h = 0.01 s. The trajectory of the vehicle over this time interval is depicted in Fig. 2. The following two inertial directions, corresponding to Nadir and Earth's magnetic field direction are measured by the inertial sensors on the vehicle:

$$d_1 = [0 \ 0 \ -1]^{\mathrm{T}}, \quad d_2 = [0.1 \ 0.975 \ -0.2]^{\mathrm{T}}.$$

For optical measurements, 8 beacons are placed at the corners of the room. It has been assumed that at least two beacons are measured at each time instant. The observed directions in the body-fixed frame are simulated with the help of the aforementioned true states. The true quantities are disturbed by bounded, random noise with zero mean to simulate realistic measurements. Based on coarse attitude sensors like sun sensors, optical sensors and magnetometers, a random noise bounded in magnitude by 2.4° is added to the matrix $L = R^T D$ to generate measured L^m . Similarly, random noise bounded in magnitude by 0.97°/s and 0.025m/s, which are close to actual noise levels of coarse rate gyros, are added to Ω , ν to generate measured Ω^m and ν^m , respectively. The scalar gain m = 1.5 and the dissipation term is chosen to be, l = 0.1. Furthermore, the value of the gains k_p and κ are chosen to be $k_p = 150$ and $\kappa = 100$. The state estimates are initialized at:

$$\hat{g} = I$$
, $\hat{\Omega}_0 = [0.1 \ 0.45 \ 0.05]^{\mathrm{T}}$ rad/s,
and $\hat{\nu}_0 = [2.05 \ 0.64 \ 1.29]^{\mathrm{T}}$ m/s.

It is to be noted that the estimation scheme presented in Sect. 5 is a discrete-time almost global asymptotically stable estimation scheme for simultaneous estimation

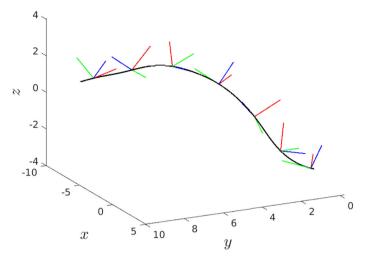
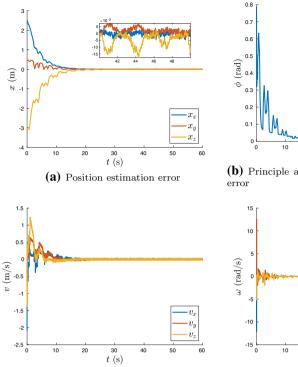


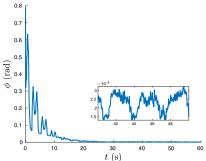
Fig. 2 Trajectory of the body

of position, attitude, linear velocity and angular velocity. The filter equations in (44) are implicit and therefore have to be solved simultaneously at each time instant in order to obtain estimates for that time instant. We start with a set of random initial state estimates as given above. With the help of the procedure mentioned in the previous paragraph, true values of inertial and optical directions, angular velocity and linear velocity are disturbed with random noise to generate realistic measurements and these measurements are fed into the filter equations at each time instant in real time. The equations in (44) are simultaneously solved with the help of fsolve available in MATLAB at each time instant to generate estimated values of state trajectories for 60s. The estimates are then compared with initially generated true state trajectories. The position and attitude estimation error are shown in Fig. 3a and b, respectively. We observe that both the position and the attitude errors converge to a bounded set around the equilibrium after about 30s. The size of the bounded set is dictated by the noise magnitudes. The corresponding velocity error plots are shown in Fig. 3c and d, respectively, and show desired performance. It is important to note that we are not required to assume the noise distribution properties.

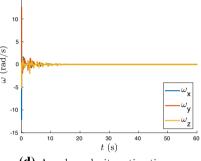
7 Conclusions

An asymptotically stable and optimal discrete-time rigid body pose estimator has been presented in this work. Suitable artificial potential energy and kinetic energy-like functions of state estimation errors were used to construct a Lagrangian in discrete time. The discrete Lagrange–d'Alembert principle was applied to this Lagrangian to obtain an optimal filtering scheme. The dissipation terms were calculated through a discrete Lyapunov analysis carried out on a Morse–Lyapunov function that corresponds to the total energy function constructed from the kinetic and potential energy-like terms used to construct the Lagrangian. The theoretical assertions are supported through realistic





(b) Principle angle of the attitude estimation error



(d) Angular velocity estimation error

Fig. 3 Performance of the estimator

(c) Translational velocity estimation error

numerical simulations. It has been observed that the estimation errors converge to a bounded neighborhood of the true states. The rates of convergence and domain of convergence can be controlled by changing scalar gains associated with the potential and kinetic energy-like terms that make up the Lagrangian. Future work in this domain would look into designing an explicit filtering scheme by constructing a suitable cost function, so that numerical computations are faster and simpler for onboard implementation.

References

- Amelin, K., Miller, A.: An algorithm for refinement of the position of a light uav on the basis of kalman filtering of bearing measurements. J. Commun. Technol. Electron. 59(6), 622–631 (2014)
- Baraff, D.: An introduction to physically based modeling: rigid body simulation i—unconstrained rigid body dynamics. SIGGRAPH course notes 82, (1997)
- Bhatt, M., Sanyal, A.K., Sukumar, S.: Optimal multi-rate rigid body attitude estimation based on lagranged'alembert principle. arXiv preprint arXiv:2008.04104 (2020)
- Bhatt, M., Sukumar, S., Sanyal, A.K.: Rigid body geometric attitude estimator using multi-rate sensors. In: 2020 59th IEEE Conference on Decision and Control (CDC), pp. 1511–1516. IEEE (2020)

Featherstone, R.: Rigid Body Dynamics Algorithms. Springer, New York (2014)

- Filipe, N., Kontitsis, M., Tsiotras, P.: Extended kalman filter for spacecraft pose estimation using dual quaternions. J. Guid. Control. Dyn. 38(9), 1625–1641 (2015)
- Izadi, M., Sanyal, A.K.: Rigid body attitude estimation based on the Lagrange-d'Alembert principle. Automatica 50(10), 2570–2577 (2014)
- Izadi, M., Sanyal, A.K.: Rigid body pose estimation based on the Lagrange-d'Alembert principle. Automatica 71, 78–88 (2016)
- La Salle, J. P.: The stability of dynamical systems. Society for Industrial and Applied Mathematics 25, (1976)
- Mahony, R., Hamel, T., Pflimlin, J.M.: Nonlinear complementary filters on the special orthogonal group. IEEE Trans. Autom. Control 53(5), 1203–1218 (2008)
- Marsden, J.E., West, M.: Discrete mechanics and variational integrators. Acta Numer. 10, 357-514 (2001)
- Rehbinder, H., Ghosh, B.K.: Pose estimation using line-based dynamic vision and inertial sensors. IEEE Trans. Autom. Control 48(2), 186–199 (2003)
- Vasconcelos, J.F., Cunha, R., Silvestre, C., Oliveira, P.: Landmark based nonlinear observer for rigid body attitude and position estimation. In: 2007 46th IEEE Conference on Decision and Control, pp. 1033– 1038. IEEE (2007)
- Vasconcelos, J.F., Cunha, R., Silvestre, C., Oliveira, P.: A nonlinear position and attitude observer on se (3) using landmark measurements. Syst. Control Lett. 59(3–4), 155–166 (2010)
- Vasconcelos, J.F., Silvestre, C., Oliveira, P.: A nonlinear gps/imu based observer for rigid body attitude and position estimation. In: 2008 47th IEEE Conference on Decision and Control, pp. 1255–1260. IEEE (2008)
- Vertechy, R., Castelli, V.P.: Accurate and fast body pose estimation by three point position data. Mech. Mach. Theory 42(9), 1170–1183 (2007)
- Wahba, G.: A least squares estimate of satellite attitude. SIAM Rev. 7(3), 409 (1965)
- Zamani, M., Trumpf, J., Mahony, R.: Minimum-energy filtering for attitude estimation. IEEE Trans. Autom. Control 58(11), 2917–2921 (2013)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Springer Nature or its licensor holds exclusive rights to this article under a publishing agreement with the author(s) or other rightsholder(s); author self-archiving of the accepted manuscript version of this article is solely governed by the terms of such publishing agreement and applicable law.